

Non-thermal heating of plasma particles due to relativistic nonlinearity of lasers

= Modeling cosmic-ray acceleration in laboratory =

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1. Laser acceleration by wake-field induced by ponderomotive force
2. An experiment shows no mono but power-law heating (Pinch to Chance)
3. Stochasticity of the particle motion by filamentation and reflection of laser
4. Power law spectrum by nonlocal diffusion in p-space. (FFPE)
5. Generation of waves and cosmic-rays in relativistic collisionless shocks in Universe
6. Proposal to NIFS

核融合科学研究所ユニット構築会議への提案として用意した。ZOOM会議。2021年7月5日

Main reference: HT book-1; <https://www.springer.com/gp/book/9783030496128>

Laser wake-field acceleration

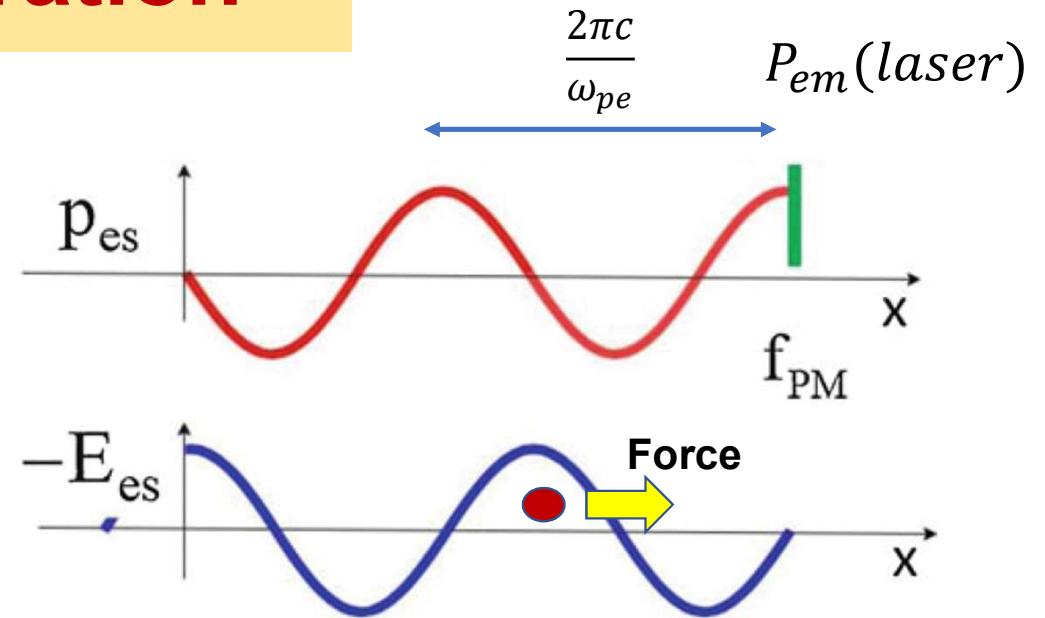
$$\frac{d^2 p_{es}}{dt^2} = -\frac{\omega_{p0}^2}{\gamma} p_{es} + p_{em} \delta(t - x_0/c)$$

$$E = \frac{e}{\epsilon_0} n_e \frac{c}{\omega_{pe}} \sim 100 \text{ [GV/m]}$$

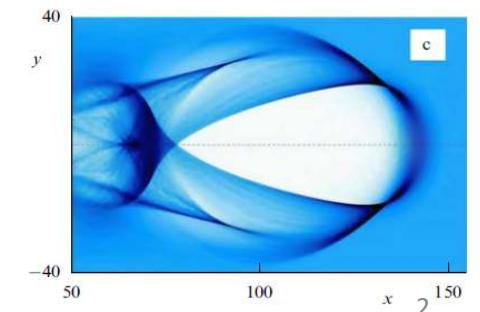
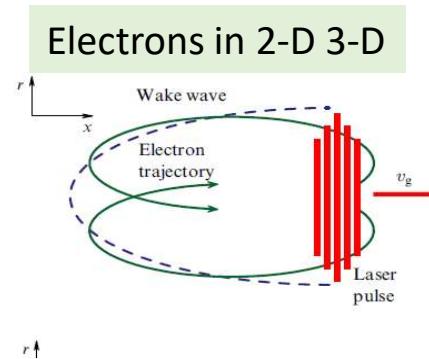
$$E = 10 \text{ [MV/m]} \text{ (LINAC)}$$



XFEL, DESY, Hamburg (17GeV/1.7km)



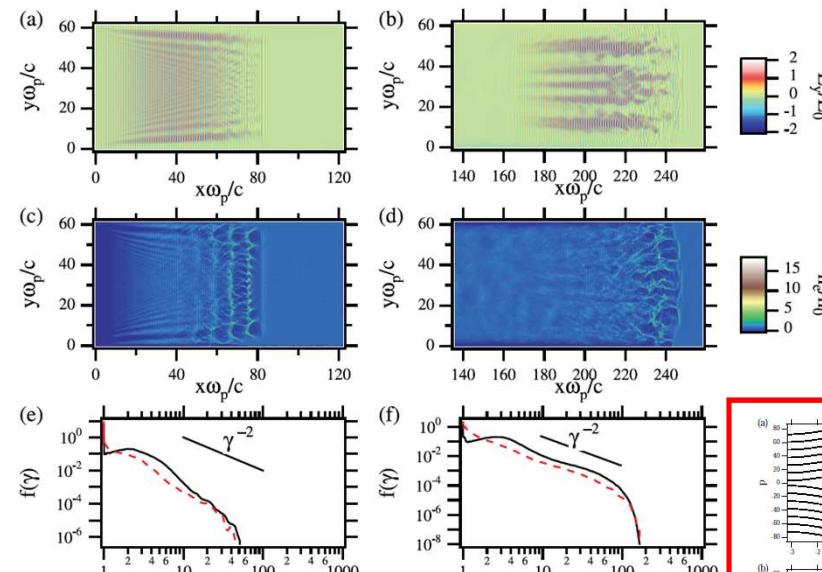
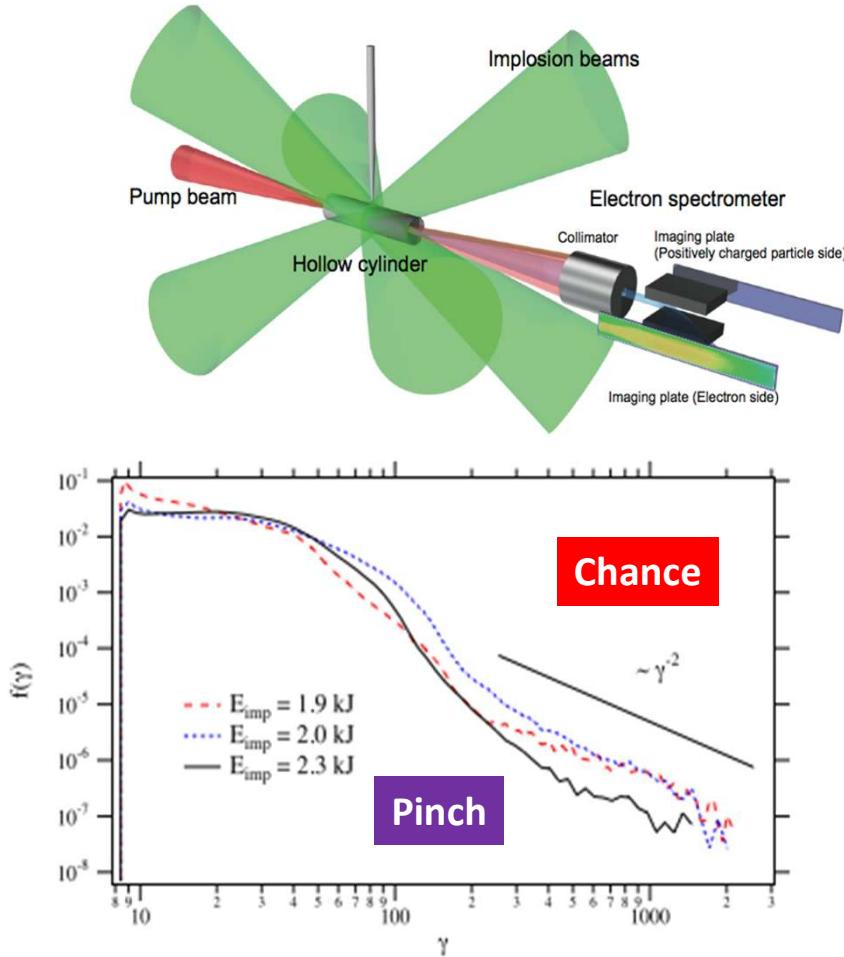
HT book-1 p. 206



LFEX experiment at ILE

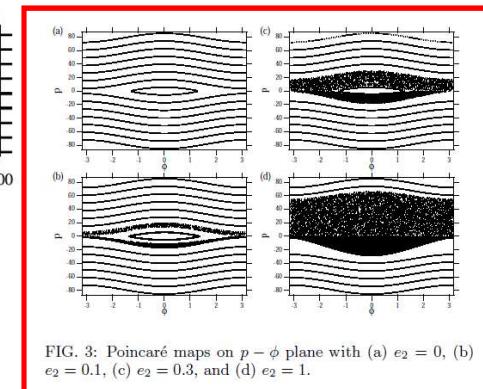


蔵満 (阪大)

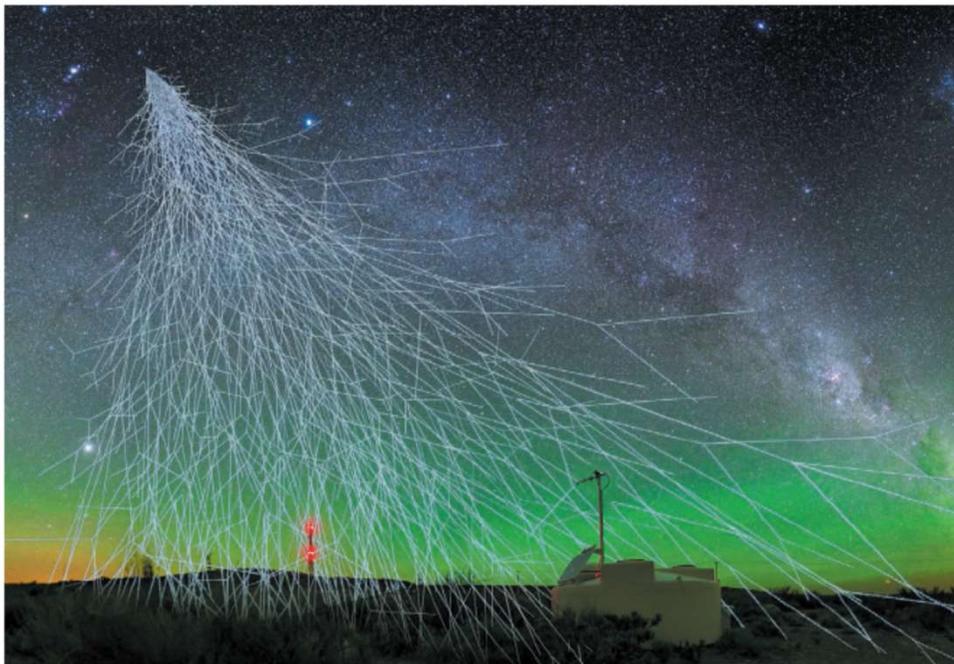


Y. Kuramitsu, *Model Experiment of Cosmic Ray Acceleration due to an Incoherent Wakefield Induced by an Intensive Laser Pulse*, Physics of Plasmas (Letter), **18**, 010201 (4 pages), (2011),

Y. Kuramitsu et al, *NONTHERMAL ACCELERATION OF CHARGED PARTICLES DUE TO AN INCOHERENT WAKEFIELD INDUCED BY A LARGE-AMPLITUDE LIGHT PULSE*, The Astrophysical Journal, 682: L113 (2008)

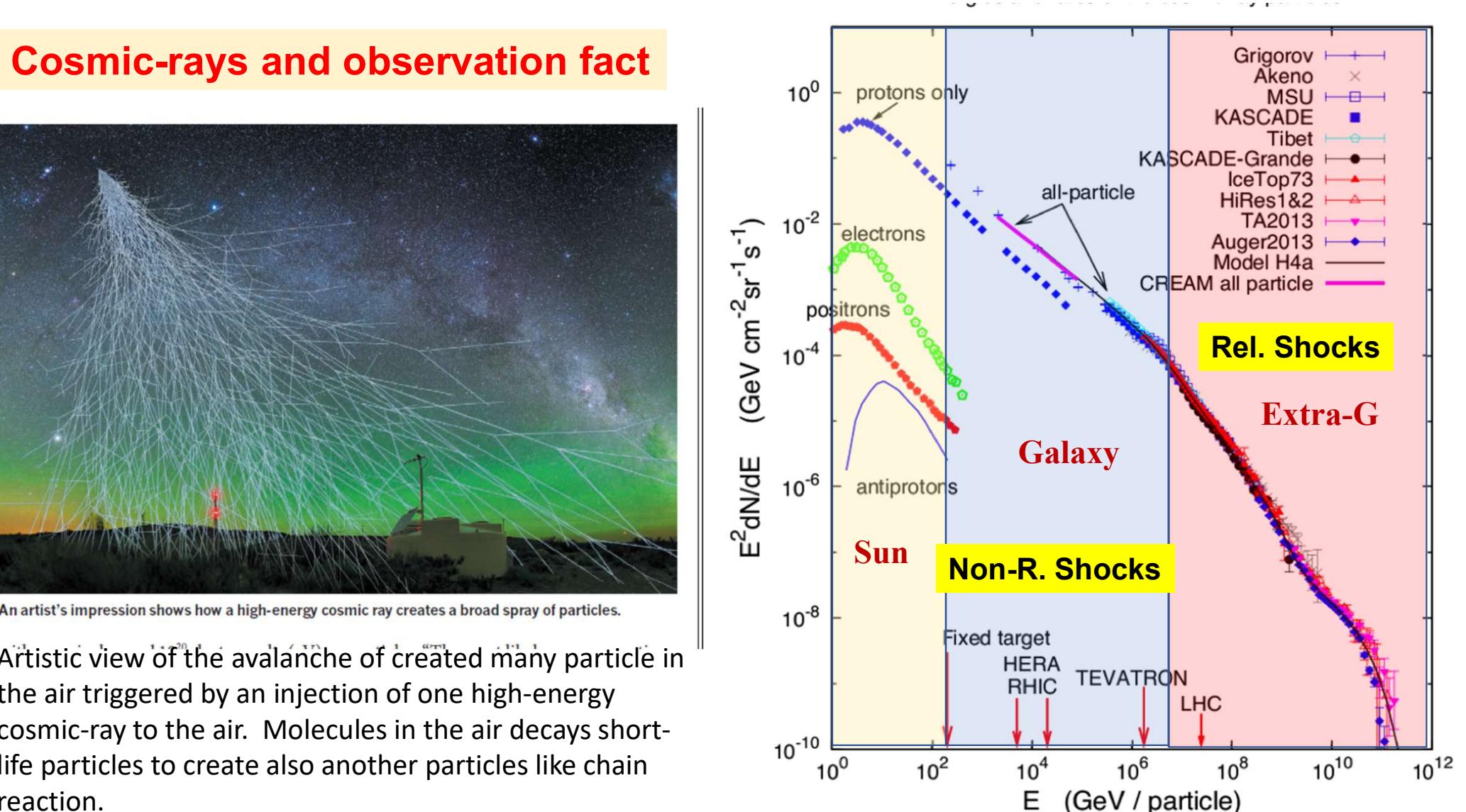


Cosmic-rays and observation fact



Artistic view of the avalanche of created many particle in the air triggered by an injection of one high-energy cosmic-ray to the air. Molecules in the air decays short-life particles to create also another particles like chain reaction.

D. Castelvecchi, Nature **549**, 440 (2017).



T. Gaisser, arXiv:1704.00788v1, astro-ph.HE (2017)

Standard model of accelerating cosmic-rays (DSA)

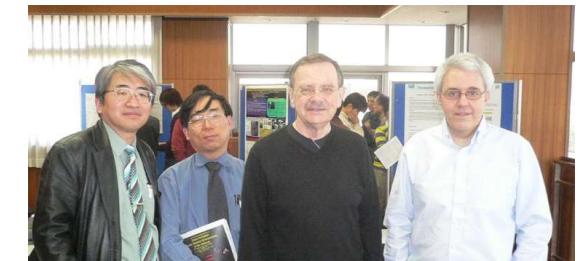
=Non-relativistic shocks=

40 years old

E. Fermi, Astrophys. J. **119**, 1 (1954).

A. R. Bell, Mon. Not. R. Astro. Soc. **182**, 147 (1978).

R. D. Blandford & J. P. Ostriker, Astrophysical J., **221**, L29 (1978)

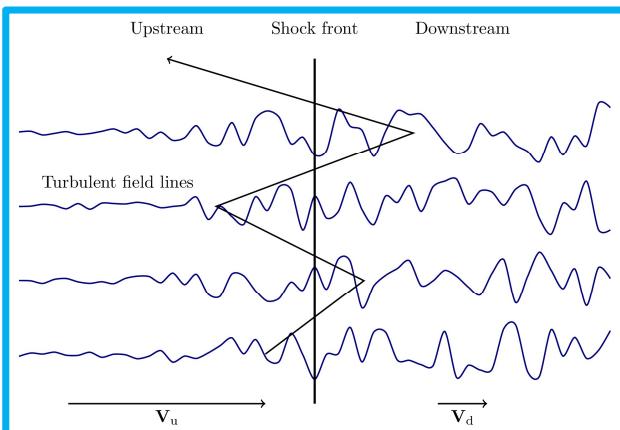


A. R. Bell

Tokyo I T, Feb. 22, 2007

Power law for non-thermal particles

Collisionless shock



$$E_n = E_0(1 + \varepsilon)^n$$

$$N(E) \propto N_0(1 - P)^n$$

$$\varepsilon = \Delta E / E$$

$$\frac{dN(E)}{dE} \propto E^{-\alpha}$$

$$\alpha = 1 - \ln(1 - P) / \ln(1 + \varepsilon) \approx 1 + P / \varepsilon$$

$$\alpha = 1 + P / \varepsilon = 1 + \frac{3}{R - 1}$$

$N(E) \sim E^{-2}$ for $R=4$

Relativistic shocks?

X-ray image of SN1006 and DSA

1. Relativistic Synchrotron Emission

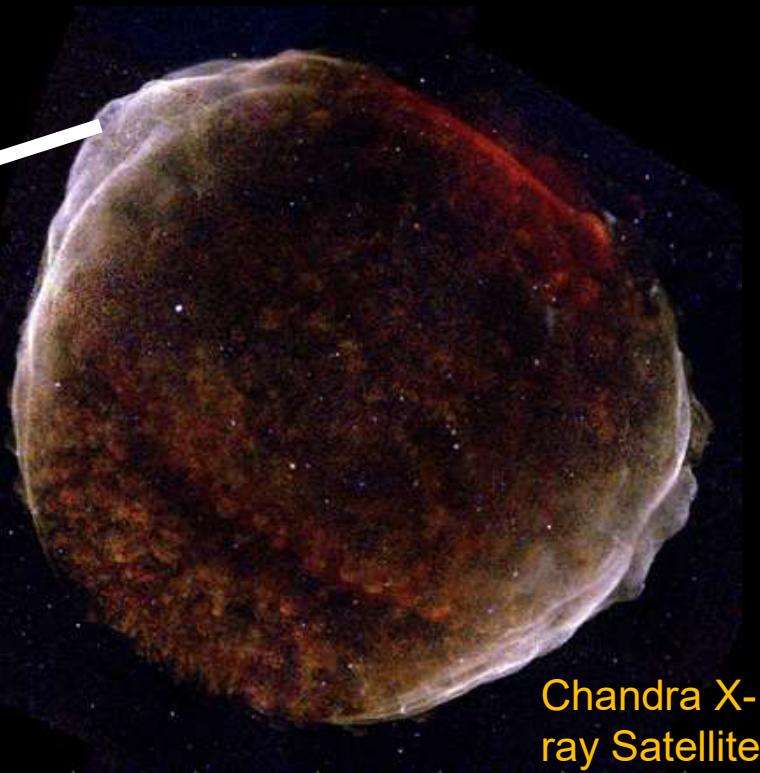
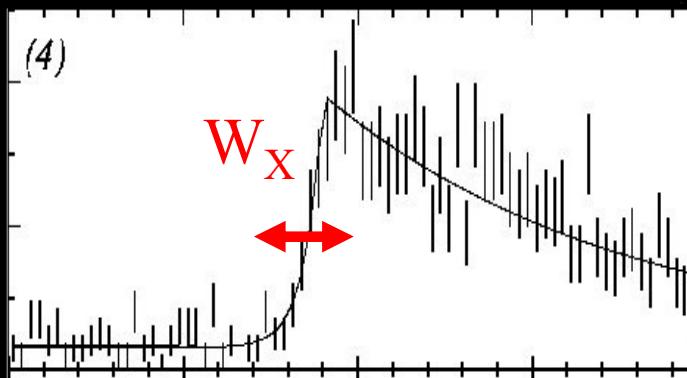
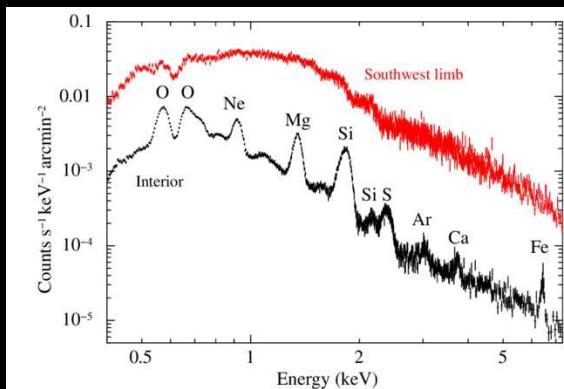
$$h\nu = 2keV \left(\frac{B}{10\mu G} \right) \left(\frac{E_e}{10^{14}eV} \right)^2$$

SN1006

$u_0 = 4330 \text{ km/s}$
diameter = 60 ly

2. Shock Thickness (Observation)

$$W_X = 1 \times 10^{17} \text{ cm} (=1/400 l_{\text{mfp}})$$



G. Cassam-Chenai et al., *Astrophys. J.* **680**, 1180 (2008)

A. Bamba et al., *Astrophys. J.* **621**, 793 (2005).

Chandra X-ray Satellite

**We have demonstrated non-relativistic collisionless
shock formation in experiment with NIF laser**

**John Dawson Award for Excellence in Plasma Physics Research
American Physical Society, 2020 Awardee**

<https://www.aps.org/programs/honors/prizes/dawson.cfm>

**Nonlinear Physics of Magnetic Turbulence, Shock Formation, and Particle
Acceleration via Weibel Instability in Laser Astrophysics**

= Theory, PIC Simulation and Laser Experiment =

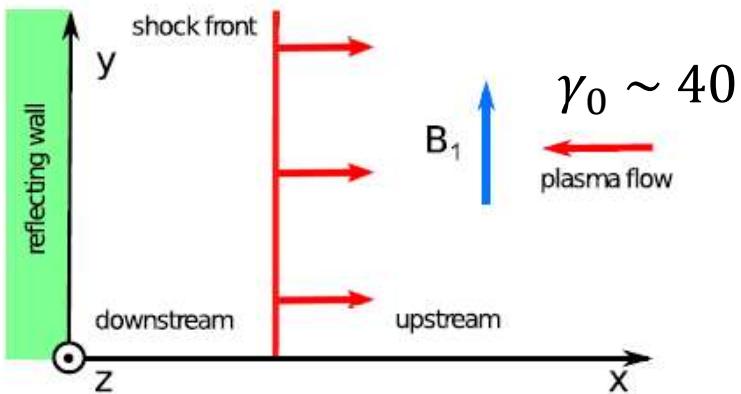
Presented at Prof. Yoshida's final seminar, March 7-8, 2021 by ZOOM
(Video of my talk)

<http://www.ppl.k.u-tokyo.ac.jp/takabe.mp4>

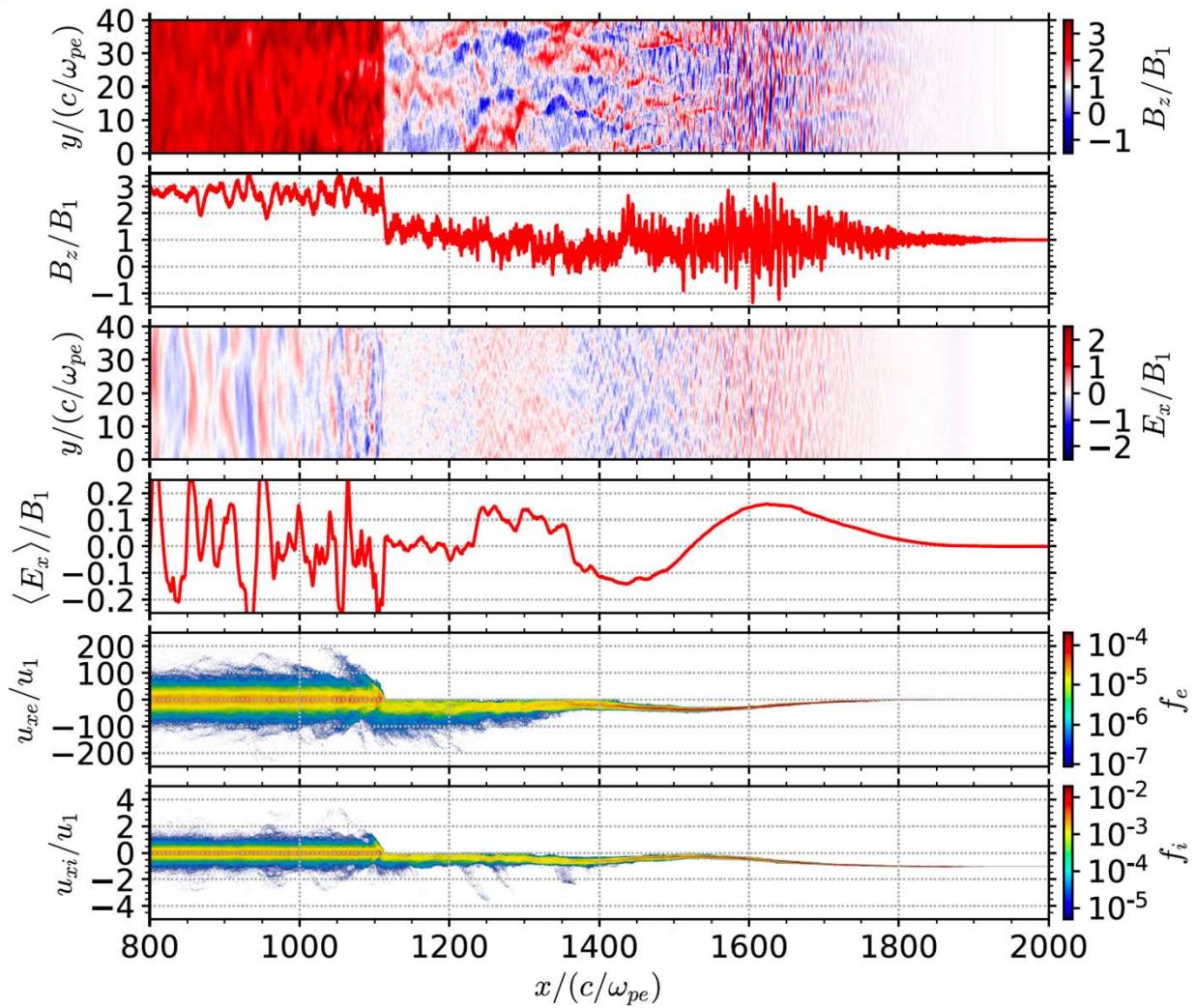
Next is relativistic shocks and cosmic-ray acceleration physics

Global Structure of 2D ion- e^- relativistic shocks

Previous one-dimensional (1D) particle-in-cell (PIC) simulations showed that **synchrotron maser instability (SMI)** is the significant dissipation mechanism for relativistic magnetized shocks (e.g., Langdon et al. 1988; Gallant et al. 1992; Hoshino et al. 1992; Amato & Arons 2006).

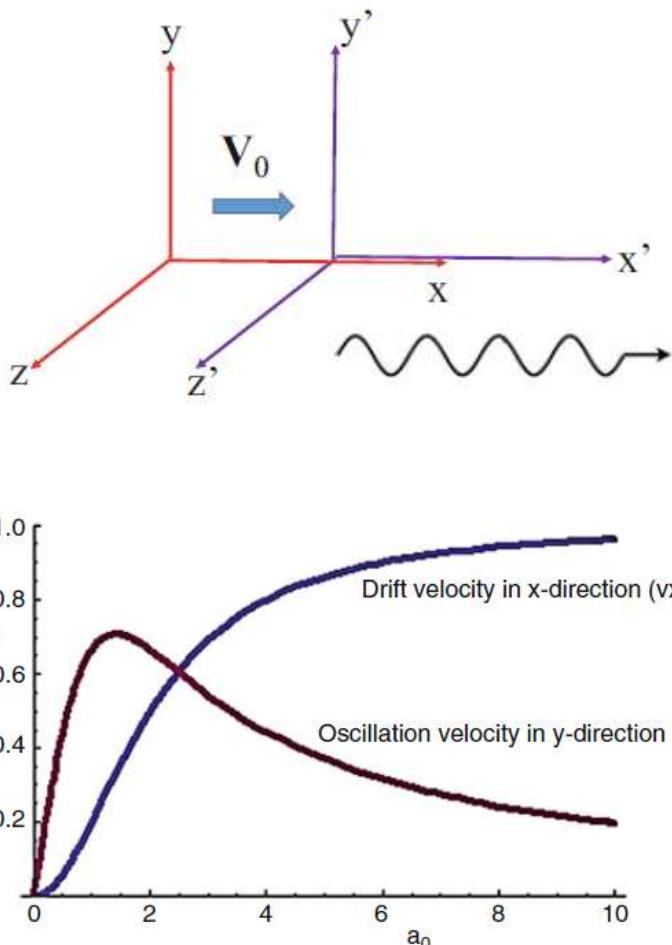


M. Iwamoto et al., *Precursor Wave Amplification by Ion-Electron Coupling through Wakefield in Relativistic Shocks*, The Astrophysical Journal Letters, 883:L35, (2019).



Relativistic EM waves (Lasers)

= Relativistic shocks =



$$\mathbf{E}'_{\perp} = \mathbf{E}_{\parallel} + \gamma_0 (\mathbf{E}_{\perp} + \mathbf{V}_0 \times \mathbf{B}_{\perp})$$

$$E'_{\perp} = \gamma_0 (1 - \beta_0) E_{\perp} \quad \gamma_0 (1 - \beta_0) = \sqrt{\frac{1 - \beta_0}{1 + \beta_0}}$$

$$(\omega', k') = \sqrt{\frac{1 - \beta_0}{1 + \beta_0}} (\omega, k)$$

$$A = A_0 \cos(kx - \omega t)$$

$$a_0 = \frac{eA_0}{mc} = \frac{eE_0}{mc\omega} = \frac{v_{os}}{c}$$

Strength parameter

Lorentz invariant

$$a_0 = 0.85 \sqrt{I_{18} \lambda_{\mu m}^2}$$

HT book-1, Chapter 5

An electron motion in relativistic laser (Nonlinear)

$$\frac{dp_x}{dt} = -ev_y B_z - eE_x$$

$$\frac{dp_y}{dt} = -eE_y + ev_x B_z$$

$$mc^2 \frac{d\gamma}{dt} = -ev_y E_y - ev_x E_x$$

$$E_y = -\frac{\partial A}{\partial t} \quad A = A_0 \cos(\xi)$$

$$B_z = \frac{\partial A}{\partial z} \quad \xi = \omega(t - x/c) + \xi_0$$

$$\frac{dp_x}{dt} = -\beta_y \frac{\partial a}{\partial x}$$

$$\frac{dp_y}{dt} = \frac{da}{dt}$$

$$\frac{d\gamma}{dt} = \beta_y \frac{\partial a}{\partial t}$$

$$p_y - a = \beta \quad \gamma - p_x = \alpha$$

$$p_x = \frac{1}{2\alpha} p_y^2 + \frac{1-\alpha^2}{2\alpha}$$

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Solution of an electron motion



Incident laser

Non-linear oscillation

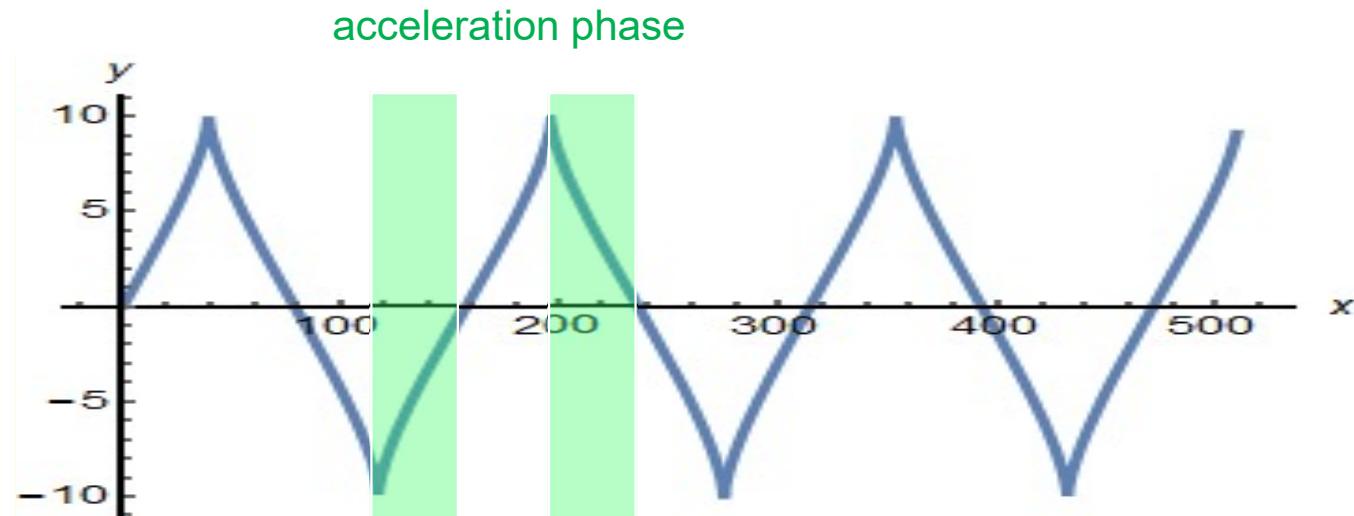
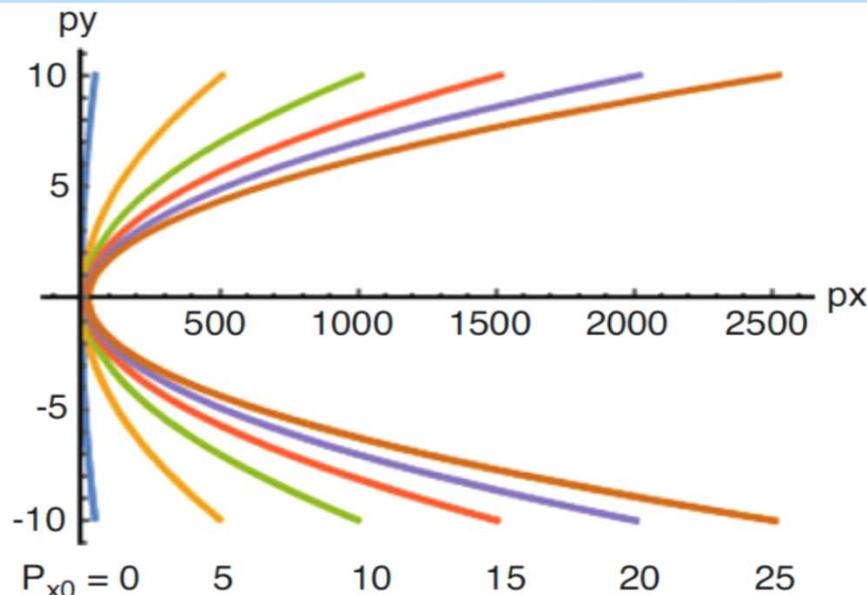
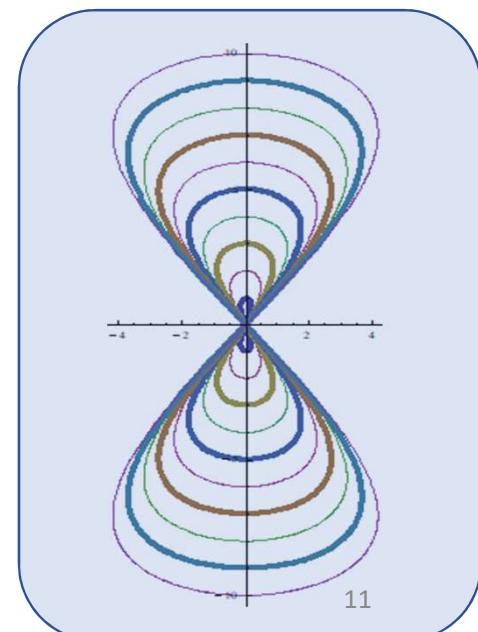
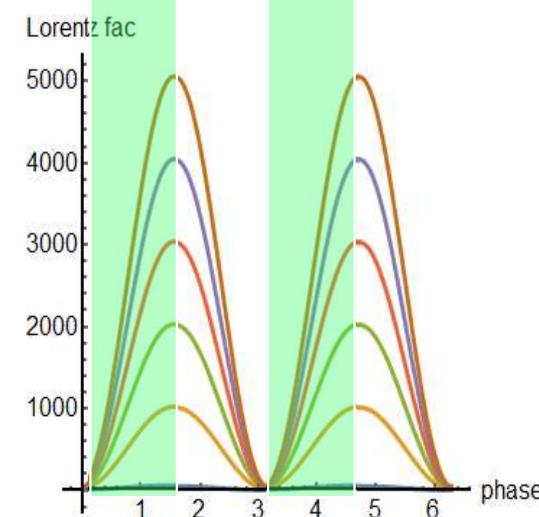


Figure of 8 motion



Motion is chaotic due to reflected laser and filamentation

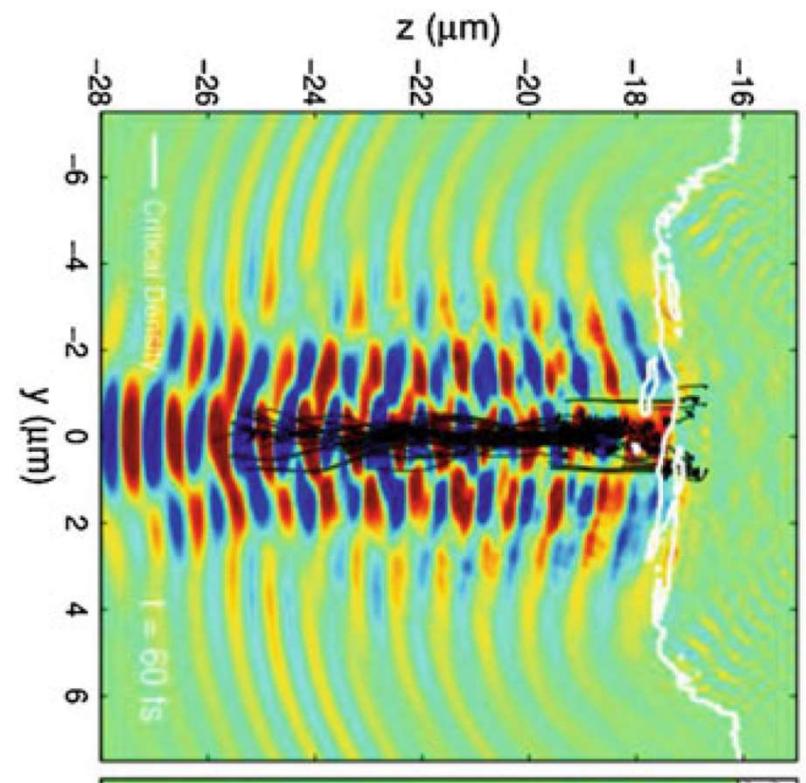
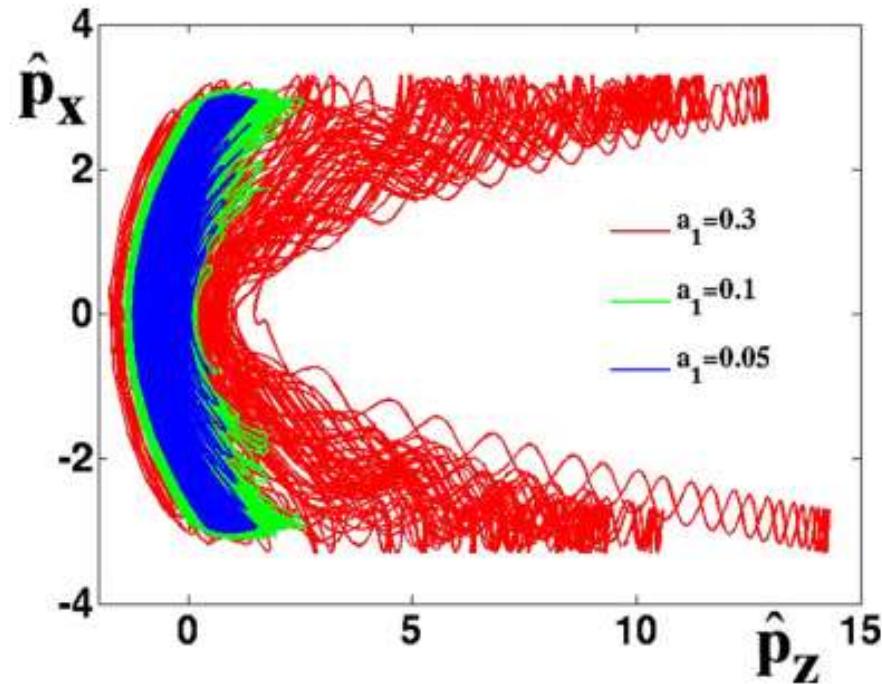


FIG. 6. $a = 3$, $P_\perp = 0$, $\gamma_0 = \sqrt{1 + a^2}$ (initial value of γ). Coordinates of the mechanical momentum of a single charged particle for three values of a_1 .

S. Rassou, A. Bourdier, and M. Drouin, Phys. Plasmas 21, 083101(2014)

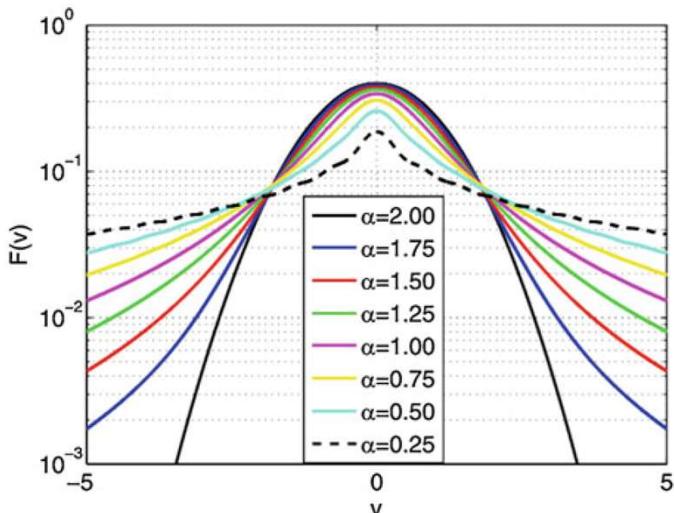
HT book-1, Chapter 9

Levy's nonlocal transport (FFPE) can explain the non-thermal components

$$\frac{\partial}{\partial t} f(p, t) = D \frac{\partial^2}{\partial p^2} f(p, t), \quad D = \left\langle \frac{(\Delta p)^2}{2\Delta t} \right\rangle$$

$$\frac{\partial^\alpha}{\partial |p|^\alpha} \equiv \frac{1}{2\pi} \int_{-\infty}^{\infty} |k|^\alpha e^{-ikp} dk \quad -k^2 \quad \Rightarrow \quad -|k|^\alpha$$

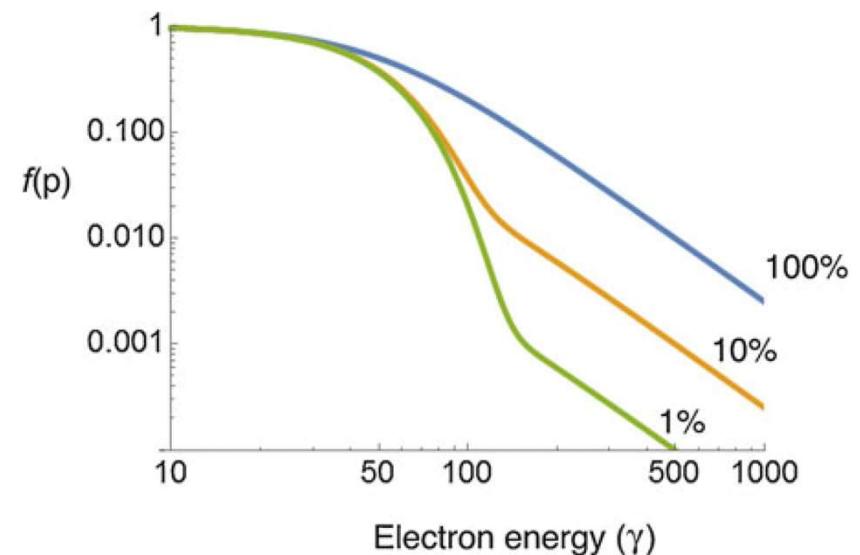
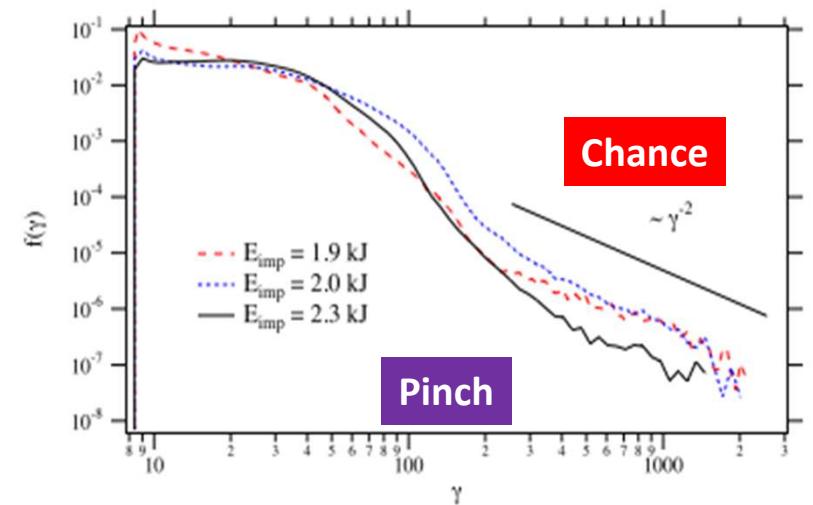
$$\frac{\partial}{\partial t} f = D \frac{\partial^\alpha}{\partial |p|^\alpha} f \quad f^\alpha(p, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp(-Dt|k|^\alpha + ikp) dk$$



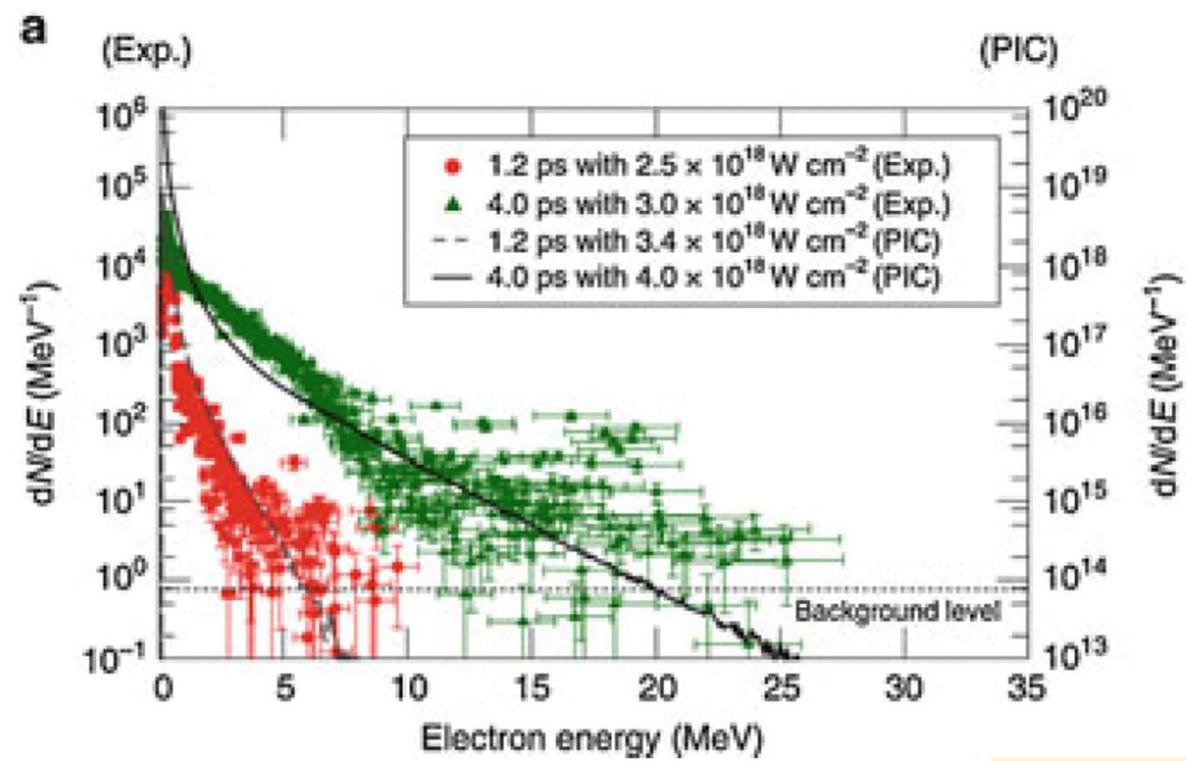
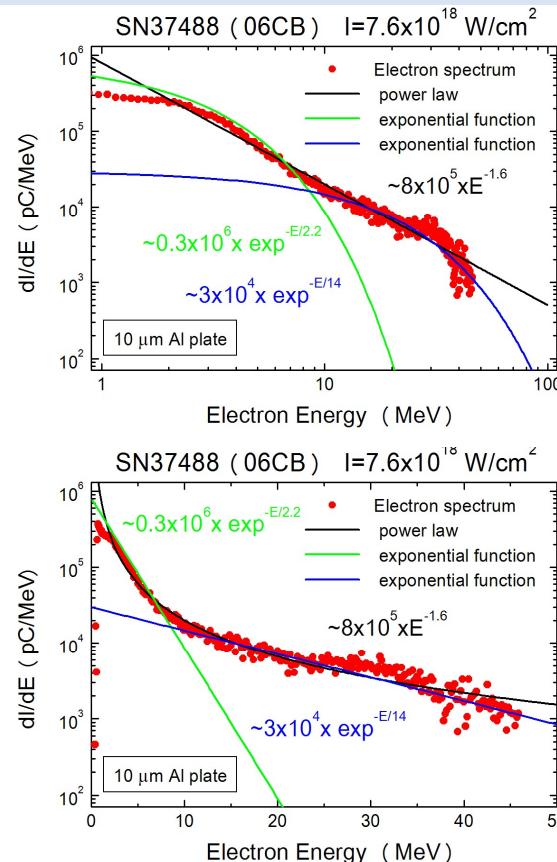
C. Tsallis et al., Phys. Rev. Lett. 75, 3589 (1995)

HT book-1, Chapter 9

J. Anderson et al., Phys. Plasmas 21, 122109 (2014)



LFEX experiment (ILE) can be explained by chaotic heating with relativistic intensity lasers



HT book-1, p. 353

S. Kojima, et al. Electromagnetic field growth triggering super-ponderomotive electron acceleration during multi-picosecond laser-plasma interaction. Commun Phys 2, 99 (2019).

EM waves induced by collisionless shocks in Space

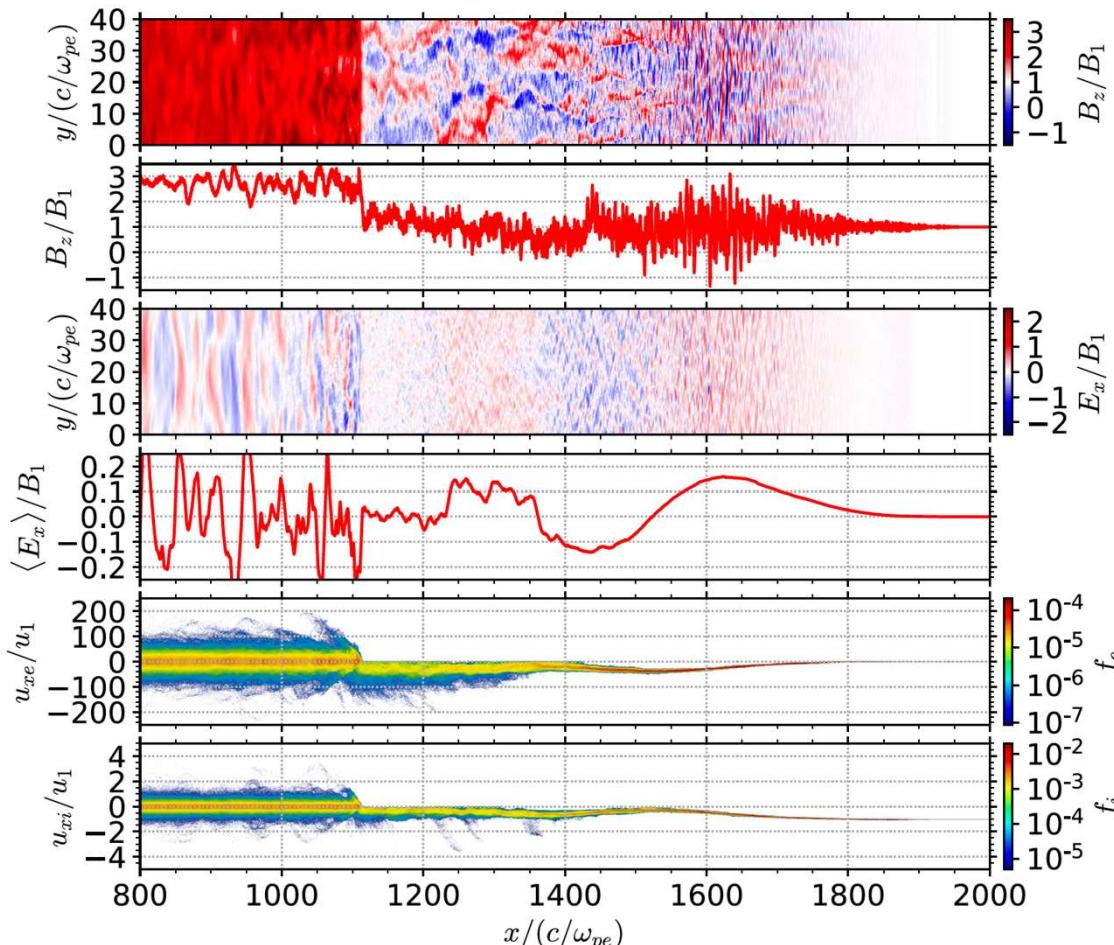


Figure 1. Shock structure for $\sigma_e = 5$ at $\omega_{pe} = 2000$. The out-of-plane magnetic field B_z , the 1D profile of B_z , the longitudinal electric field E_x , the y -averaged electric field E_x , and the phase space plots in the x - u_x plane for electrons and ions are shown.

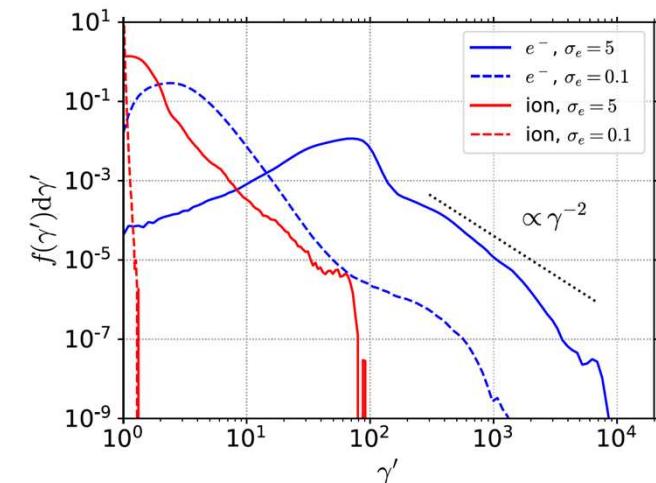


Figure 4. Near-upstream energy spectra of electrons (blue) and ions (red) for $\sigma_e = 5$ (solid lines) and $\sigma_e = 0.1$ (dashed lines) measured in the proper frame.

M. Iwamoto et al., *Precursor Wave Amplification by Ion–Electron Coupling through Wakefield in Relativistic Shocks*, The Astrophysical Journal Letters, 883:L35, (2019).

HT book-1, p. 354

Proposal to the NIFS unit

- Study the theory of chaotic acceleration of electrons by relativistic EM waves
- Relate the theory to the physical mechanism of cosmic-ray acceleration in relativistic regime ($E > 10^{15}$ eV)
- Carry out big computing to visualize the physical mechanism
- Design a verification experiment with reasonable facility parameters
- Perform the experiment with any of international facilities like ELIs in EU.
- Challenge to propose a new theoretical mode and experimental proof to the big community of the cosmic-ray research.